# ECED 3300 <br> Tutorial 5 

## Problem 1

A point charge $Q$ is located at the origin. Find the electric flux that crosses
(a) a sphere of radius $R$ centered at the origin;
(b) the annulus, described by the condition $\alpha_{1} \leq \phi \leq \alpha_{2}$, which lies on a sphere of radius $R$ centered at the origin;
(c) the patch, specified by $\alpha_{1} \leq \phi \leq \alpha_{2}$ and $\theta_{1} \leq \theta \leq \theta_{2}$, which lies on a sphere of radius $R$ centered at the origin.

## Solution

(a) By definition, the flux through the entire sphere is equal to the total charge enclosed by the surface of the sphere, $\Psi=Q$.
(b) By definition, $\Psi=\int_{S} \mathbf{D} \cdot d \mathbf{S}$. In our case, $d \mathbf{S}=\mathbf{a}_{r} R^{2} \sin \theta d \theta d \phi$, and $\mathbf{D}=Q /\left(4 \pi r^{2}\right) \mathbf{a}_{r}$ (the flux density due to a point charge at the origin). It follows that

$$
\Psi=\frac{Q}{4 \pi R^{2}} R^{2} \int_{0}^{\pi} \sin \theta d \theta \int_{\alpha_{1}}^{\alpha_{2}} d \phi=\frac{Q}{4 \pi}\left(\alpha_{2}-\alpha_{1}\right) \times\left.\cos \theta\right|_{\pi} ^{0}=\frac{Q\left(\alpha_{2}-\alpha_{1}\right)}{2 \pi}
$$

(c) Similarly to part (b):

$$
\Psi=\frac{Q}{4 \pi} \int_{\theta_{1}}^{\theta_{2}} \sin \theta d \theta \int_{\alpha_{1}}^{\alpha_{2}} d \phi=\frac{Q\left(\cos \theta_{1}-\cos \theta_{2}\right)\left(\alpha_{2}-\alpha_{1}\right)}{4 \pi} .
$$

## Problem 2

The interior of an infinitely long cylinder of radius $a$ is filled with an axially symmetric volume charge distribution, $\rho_{v}(\rho)=\kappa \frac{\rho}{a},\left(C / m^{3}\right)$, where $\kappa=$ const. The charged cylinder is surrounded by an infinitely long, thin cylindrical shell of the radius $b, b>a$, which carries the charge $\rho_{l},(C / m)$, per unit length. Assuming the dielectric constants of all media to be equal to unity, find the electric field everywhere.

## Solution

By axial symmetry of the system, the electric field can only have a radial component; a cylinder of length $L$ and radius $r$ will be chosen as a special Gaussian surface. We can then apply Gauss's law to the charge distribution enclosed by such a surface. Let us consider the following cases

1. $\rho \leq a$. Consider a cylinder of radius $\rho, \rho \leq a$ and length $L$ as a Gaussian surface. Applying Gauss's law, we obtain the equation

$$
D 2 \pi \rho L=L \int_{0}^{2 \pi} d \phi \int_{0}^{\rho} \rho^{\prime} d \rho^{\prime} \kappa \rho^{\prime} / a
$$

Performing the elementary integration on the right-hand side, we arrive at

$$
D 2 \pi \rho L=2 \pi L \frac{\kappa \rho^{3}}{3 a} \Longrightarrow D=\frac{\kappa \rho^{2}}{3 a}
$$

Due to the relation between the flux density and the electric field, $\mathbf{D}=\epsilon_{0} \mathbf{E}$, we obtain at once the final expression for the field

$$
\mathbf{E}=\frac{\kappa \rho^{2}}{3 \epsilon_{0} a} \mathbf{a}_{\rho}
$$

2. $a \leq \rho \leq b$. By the same token,

$$
D 2 \pi \rho L=2 \pi L \int_{0}^{a} \rho^{\prime} d \rho^{\prime} \kappa \rho^{\prime} / a
$$

It follows that

$$
\mathbf{E}=\frac{\kappa a^{2}}{3 \epsilon_{0} \rho} \mathbf{a}_{\rho}
$$

in this region.
3. $\rho \geq b$. Here, the flux density is determined by the total charge inside the cylinder and the shell:

$$
D 2 \pi \rho L=2 \pi L \frac{\kappa a^{2}}{3}+\rho_{l} L
$$

Thus

$$
\mathbf{E}=\frac{\kappa a^{2}}{3 \epsilon_{0} \rho}\left(1+\frac{3 \rho_{l}}{2 \pi \kappa a^{2}}\right) \mathbf{a}_{\rho}
$$

## Problem 3

Determine the charge density due to the following electric charge distributions
(a) $\mathbf{D}=4 \rho \sin \phi \mathbf{a}_{\rho}+2 \rho \cos \phi \mathbf{a}_{\phi}+2 z^{2} \mathbf{a}_{z}, C / m^{2}$.
(b) $\mathbf{D}=\frac{2 \cos \theta}{r^{3}} \mathbf{a}_{r}+\frac{\sin \theta}{r^{3}} \mathbf{a}_{\theta}, C / m^{2}$.

## Solution

Using the differential form of Gauss's law,

$$
\nabla \cdot \mathbf{D}=\rho_{v}
$$

(a) Thus

$$
\rho_{v}=\frac{1}{\rho} \frac{\partial}{\partial \rho}\left(\rho D_{\rho}\right)+\frac{1}{\rho} \frac{\partial D_{\phi}}{\partial \phi}+\frac{\partial D_{z}}{\partial z} .
$$

It follows that

$$
\rho_{v}=\frac{1}{\rho} \frac{\partial}{\partial \rho}\left(4 \rho^{2} \sin \phi\right)+2 \frac{\partial(\cos \phi)}{\partial \phi}+\frac{\partial\left(2 z^{2}\right)}{\partial z}=6 \sin \phi+4 z .
$$

(b) By the same token,

$$
\rho_{v}=\frac{1}{r^{2}} \frac{\partial}{\partial r}\left(r^{2} D_{r}\right)+\frac{1}{r \sin \theta} \frac{\partial}{\partial \theta}\left(\sin \theta D_{\theta}\right)+\frac{1}{r \sin \theta} \frac{\partial D_{\phi}}{\partial \phi} .
$$

It can be inferred that

$$
\rho_{v}=\frac{2 \cos \theta}{r^{2}} \frac{\partial}{\partial r}\left(\frac{1}{r}\right)+\frac{1}{r^{4} \sin \theta} \frac{\partial}{\partial \theta}\left(\sin ^{2} \theta\right)=-\frac{2 \cos \theta}{r^{4}}+\frac{2 \cos \theta}{r^{4}}=0 .
$$

